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Prof. Dr. Urs Lang	Solution 9	FS 2025

9.1. Almost complex structure

An almost complex structure J on a manifold M^m is a (1, 1)-tensor field with the following property: if for every $p \in M$ we denote by $J_p: T_pM \to T_pM$ the linear map associated with J (recall Theorem T.3), then

$$J_p \circ J_p = -\mathrm{id}_{T_p M}.$$

Prove that every complex manifold admits an almost complex structure.

Hint: Composed with the differential of a complex chart $\varphi : U \to \varphi(U) \subset \mathbb{C}^n$, J_p amounts to the multiplication by *i*.

Solution. Let $\varphi: U \to \varphi(U) \subset \mathbb{C}^n$ be a chart with coordinates $(x_1, y_1, \ldots, x_n, y_n)$. As suggested by the hint, we define J locally by

$$J\left(\frac{\partial}{\partial x_i}\right) = \frac{\partial}{\partial y_i}, \quad J\left(\frac{\partial}{\partial y_i}\right) = -\frac{\partial}{\partial x_i}, \quad \text{for } i = 1, \dots, n.$$

It follows that $J_p \circ J_p = -\mathrm{id}_{T_pM}$. It remains to show that J is (globally) well-defined. Let $\psi: V \to \psi(V) \subset \mathbb{C}^n$ be another complex chart on M such that $U \cap V \neq \emptyset$, and denote the coordinates on V by $(u_1, v_1, \ldots, u_n, v_n)$, then

$$\frac{\partial}{\partial x_k} = \sum_i \frac{\partial u_i}{\partial x_k} \frac{\partial}{\partial u_i} + \frac{\partial v_i}{\partial x_k} \frac{\partial}{\partial v_i}, \quad \frac{\partial}{\partial y_k} = \sum_i \frac{\partial u_i}{\partial y_k} \frac{\partial}{\partial u_i} + \frac{\partial v_i}{\partial y_k} \frac{\partial}{\partial v_i}.$$

Since $\psi \circ \varphi^{-1}$ is biholomorphic, the Cauchy-Riemann equations imply

$$\frac{\partial u_i}{\partial x_k} = \frac{\partial v_i}{\partial y_k}, \quad \frac{\partial u_i}{\partial y_k} = -\frac{\partial v_i}{\partial x_k}.$$

Denote by J' the corresponding map defined on V with respect to ψ , then

$$J'\left(\frac{\partial}{\partial x_k}\right) = J'\left(\sum_i \frac{\partial u_i}{\partial x_k} \frac{\partial}{\partial u_i} + \frac{\partial v_i}{\partial x_k} \frac{\partial}{\partial v_i}\right)$$
$$= \sum_i \frac{\partial u_i}{\partial x_k} \frac{\partial}{\partial v_i} - \frac{\partial v_i}{\partial x_k} \frac{\partial}{\partial u_i}$$
$$= \sum_i \frac{\partial v_i}{\partial y_k} \frac{\partial}{\partial v_i} + \frac{\partial u_i}{\partial y_k} \frac{\partial}{\partial u_i}$$
$$= \frac{\partial}{\partial y_k}$$

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and similarly $J'\left(\frac{\partial}{\partial y_k}\right) = -\frac{\partial}{\partial x_k}$. This shows that J and J' coincide on $U \cap V$. \Box

9.2. Kähler manifolds

Let M be a complex manifold with an almost complex structure $J \in \Gamma(T_{1,1}M)$ (as in Exercise 1). Suppose that M is endowed with an hermitian metric, that is, $g_p(J_pv, J_pw) = g_p(v, w)$ for all $p \in M$ and $v, w \in T_pM$. Show that

$$\omega(X,Y) \coloneqq g(X,JY)$$

defines a 2-form $\omega \in \Omega^2(M)$, which is closed if and only if J is parallel (i.e. $DJ = D^{1,1}J \equiv 0$).

Solution. By definition ω is a (0, 2)-tensor field in $\Gamma(T_{0,2}M)$. We still have to show that it is antisymmetric and for that we will use that $J_p^2 = -\mathrm{id}M_p$. For $X, Y \in \Gamma(TM)$ $\omega(X,Y) = g(X,JY) = g(JX,J^2Y) = -g(JX,Y) = -g(Y,JX) = -\omega(Y,X)$, thus $\omega \in \Omega^2(M)$. In order to prove the second statement, we'll prove the following two identities

$$d\omega(X, Y, Z) = g(X, (D_Z J)Y) + g(Y, (D_X J)Z) + g(Z, (D_Y J)X)$$

$$2g(D_X(JY), Z) = d\omega(X, JY, JZ) - d\omega(X, Y, Z)$$

Let X, Y, Z, JX, JY, JZ be coordinate vector fields on a chart of M, in particular they commute. Then (see Theorem 11.3 of Differential Geometry I)

$$d\omega(X, Y, Z) = X\omega(Y, Z) - Y\omega(X, Z) + Z\omega(X, Y)$$

= $Xg(Y, JZ) - Yg(X, JZ) + Zg(X, JY)$
= $Xg(Y, JZ) + Yg(Z, JX) + Zg(X, JY)$

Thus by the compatibility of the Levi-Civita connection with g and the product rule $D_X(JZ) = (D_XJ)Z + J(D_XZ)$ for tensor derivations (similarly for the other tuples), we compute

$$\begin{split} d\omega(X,Y,Z) =& Xg(Y,JZ) + Yg(Z,JX) + Zg(X,JY) \\ =& g\left(D_XY,JZ\right) + g\left(Y,D_X(JZ)\right) + g\left(D_YZ,JX\right) \\ &+ g\left(Z,D_Y(JX)\right) + g\left(D_ZX,JY\right) + g\left(X,D_Z(JY)\right) \\ =& g\left(D_XY,JZ\right) + g\left(D_YZ,JX\right) + g\left(D_ZX,JY\right) \\ &+ g\left(Y,JD_XZ\right) + g\left(Z,JD_YX\right) + g\left(X,JD_ZY\right) \\ &+ g\left(Y,(D_XJ)Z\right) + g\left(Z,(D_YJ)X\right) + g\left(X,(D_ZJ)Y\right) \\ =& g\left(Y,(D_XJ)Z\right) + g\left(Z,(D_YJ)X\right) + g\left(X,(D_ZJ)Y\right) , \end{split}$$

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where in the last step we have used that g is hermitian and the vector fields commute. This proves the first identity.

For the second identity first note that $g((D_XJ)Y,Z) = g(D_X(JY),Z) - g(JD_XY,Z) = g(D_X(JY),Z) + g(D_XY,JZ)$. Now, by the Koszul formula we have

$$2g (D_X(JY), Z) = Xg(JY, Z) + JYg(X, Z) - Zg(X, JY)$$

= $X\omega(Z, Y) - JYg(X, JJZ) - Z\omega(X, Y)$
= $-X\omega(Y, Z) - JY\omega(X, JZ) - Z\omega(X, Y)$

and

$$2g(D_XY,JZ) = Xg(Y,JZ) + Yg(X,JZ) - JZg(X,Y)$$

= $-Xg(JZ,JJY) + Y\omega(X,Z) + JZg(X,JJY)$
= $-X\omega(JZ,JY) + Y\omega(X,Z) + JZ\omega(X,JY)$

By summing the two expression we obtain the second identity.